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## Session 8: Solutions

The Ising model in 1D is  $H = -J \sum_i s_i s_{i+1} - h \sum_i s_i$

$$Z = \sum_{\{s_i\}} e^{\beta J \sum_i s_i s_{i+1} - h \sum_i s_i}$$

Let's coarse grain: we do not touch the even sites, and we sum over the odd ones

$$Z = \sum_{\{s_i\}_{i \text{ even}}} e^{\beta h \sum_{i=-\infty}^{+\infty} s_{2i}} \sum_{\{s_i\}_{i \text{ odd}}} e^{\beta h \sum_{i=-\infty}^{+\infty} s_{2i+1} + \beta J \sum_{i=-\infty}^{+\infty} s_{2i+1} (s_{2i} + s_{2i+2})}$$

I symmetrize it =

$$= \sum_{\{s_i\}_{i \text{ even}}} e^{\beta h \sum_{i=-\infty}^{+\infty} s_{2i}} \sum_{\{s_i\}_{i \text{ odd}}} \prod_{i=-\infty}^{+\infty} e^{s_{2i+1} [\beta h + \beta J (s_{2i} + s_{2i+2})]}$$

$$= \sum_{\{s_i\}_{i \text{ even}}} e^{\beta h \sum_{i=-\infty}^{+\infty} s_{2i}} \prod_{i=-\infty}^{+\infty} [2 \cosh [\beta h + \beta J (s_{2i} + s_{2i+2})]]$$

$$= \sum_{\{s_i\}_{i \text{ even}}} e^{\beta \frac{h}{2} \sum_{i=-\infty}^{+\infty} (s_{2i} + s_{2i+2})} \prod_{i=-\infty}^{+\infty} 2 \cosh [\beta h + \beta J (s_{2i} + s_{2i+2})] =$$

$$= \sum_{\{s_i\}_{i \text{ even}}} \prod_{i=-\infty}^{+\infty} 2 e^{\beta \frac{h}{2} (s_{2i} + s_{2i+2})} \cosh [\beta h + \beta J (s_{2i} + s_{2i+2})] =$$

$$= \sum_{\{s_i\}_{i \text{ even}}} \prod_{i=-\infty}^{+\infty} 2 e^{\frac{1}{2} K_1 (s_{2i} + s_{2i+2})} \cosh [K_1 + K_2 (s_{2i} + s_{2i+2})]$$

$$K_1 = \beta h \quad K_2 = \beta J$$

How does this compare with the original Ising model before coarse graining?

$$Z = \sum_{\{s_i\}} e^{\beta h \sum_i s_i + \beta J \sum_i s_i s_{i+1}} =$$

$$= \sum_{\{s_i\}} \prod_i e^{\frac{K_1}{2} (s_i + s_{i+1}) + K_2 s_i s_{i+1}}$$

If we want to write the coarse-grained expression as a new Ising model, we have to match

$$e^{\frac{K_1'}{2} (s_i + s_{i+1}) + K_2' s_i s_{i+1} + K_0'} = 2 e^{\frac{K_1}{2} (s_i + s_{i+1}) + K_0} \cdot \cosh[K_1 + K_2 \cdot (s_i + s_{i+1})]$$

there is no problem solving it

Here we have changed  $2i \rightarrow i$  since only the even sites have remained

We can match these conditions case by case

$$\left. \begin{array}{l} ++ \\ +- \\ -+ \\ -- \end{array} \right\} \begin{cases} e^{K_1' + K_2' + K_0'} = 2 e^{K_0 + K_1} \cosh[K_1 + 2K_2] \\ e^{-K_2' + K_0'} = 2 e^{K_0} \cosh[K_1] \\ e^{-K_1' + K_2' + K_0'} = 2 e^{K_0 - K_1} \cosh[K_1 - 2K_2] \end{cases}$$

From here we can start solving the equations

- divide the first by the third:

$$e^{2k_1'} = e^{2k_1} \frac{\cosh(k_1 + 2k_2)}{\cosh(k_1 - 2k_2)}$$

- multiply (1) and (3) and divide by (2)<sup>2</sup>:

$$e^{4k_2'} = \frac{\cosh(k_1 + 2k_2) \cosh(k_1 - 2k_2)}{\cosh(k_1)^2}$$

- and at last

$$e^{4k_0'} = 8 e^{k_0} \cosh(2k_2 + k_1) \cosh(2k_2 - k_1) \cosh^2(k_1)$$

Now we have the iterative map

$$e^{2k_1'} = e^{2k_1} \frac{\cosh(k_1 + 2k_2)}{\cosh(k_1 - 2k_2)}$$

$$e^{4k_2'} = \frac{\cosh(k_1 + 2k_2) \cosh(k_1 - 2k_2)}{\cosh(k_1)^2}$$

Let's now focus on  $k_1=0$  ( $h=0$  : no field)

The two equations become

$$\begin{cases} e^{2k_1'} = \pm 1 \\ e^{4k_2'} = \cosh^2(2k_2) \end{cases} \Rightarrow k_1' = 0$$

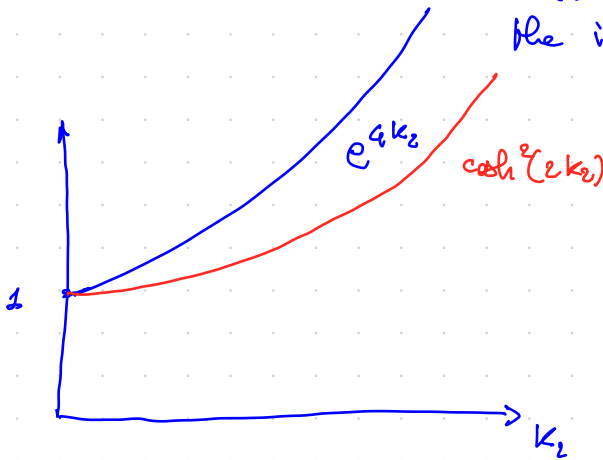
Does  $k_1' = 0$ ?  $k_1'$  multiplies  $s_i$  : its contribution to the energy depends on the state of  $s_i$ . It breaks the  $\mathbb{Z}_2$  symmetry ( $+1 \leftrightarrow -1$ )

If  $h=0$  the  $\mathbb{Z}_2$  symmetry is true, and it must be preserved by the coarse-graining  
 $\Rightarrow k_1' = 0$  is the correct answer

Let's look at the other equation:

$$e^{4k_2'} = \cosh^2(2k_2)$$

and let's solve it graphically (when solving graphically,  $k_2' = k_2$  in the graph and we look for the intersections)



Where do they intersect?

$$k_2 = 0 \quad \Rightarrow \quad \beta J = 0 \quad \Rightarrow \quad \frac{J}{k_B T} = 0 \Rightarrow T = \infty$$

$$k_2 = \infty \quad \Rightarrow \quad \beta J = \infty \quad \Rightarrow \quad \frac{J}{k_B T} = \infty \Rightarrow T = 0$$

There are no fixed points for finite  $k_2 \Rightarrow$  no critical point!

This is correct for the 1D Ising model!

What about the stability?

Let's focus on  $k_2^* = 0$  ( $T = \infty$ )

$$k_2 = k_2^* + \delta k_2 = \delta k_2$$

$$k_2' = k_2^* + \delta k_2' = \delta k_2'$$

$$e^{4\delta k_2'} = \cosh^2(2\delta k_2)$$

↓ linear expand

$$1 + 4\delta k_2' = (1 + \delta k_2)^2 = 1 + 2\delta k_2$$

$$\Rightarrow \delta k_2' = \frac{1}{2} \delta k_2$$

It is a contraction!  $k_2^* = 0$  is stable

Since the only one left is  $k_2^* = \infty$ , it must be unstable. The renormalization flow is

